

LG EXCITATION, ATTENUATION AND SOURCE SPECTRAL SCALING IN CENTRAL ASIA AND CHINA

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TABLE OF CONTENTS

I.	ABSTRACT	1
II.	STATEMENT OF WORK.....	3
III.	RESEARCH ACCOMPLISHED	4
	Data processing	4
	Agreements among 1 Hz Q measured using Lg from explosions,	
	Lg from earthquakes and Lg coda	15
	Comparison between Lg η and Lg coda η	16
	Correlation between M_0 (Lg) and ISC M_b	18
	Previously unresolved issues on the scaling between M_0 and f_c values....	19
	Scaling between M_0 and f_c values for explosions	20
	Scaling between M_0 and f_c values for earthquakes	23
	Dependence of M_0 , f_c estimates on theoretical source models	23
	Comparison between M_0 - f_c scalings for earthquakes and explosions	24
	Implications on the difference between excitations of local S and Lg.....	25
IV.	FUTURE RESEARCH	25
V.	REFERENCES.....	26
VI.	PUBLICATIONS	29
VII.	PROFESSIONAL PERSONNEL.....	29
VIII.	PAPERS PRESENTED AT MEETINGS.....	29

LIST OF TABLES AND FIGURES

Table 1	5
Table 2	6
Figure 1	7
Figure 2	8
Figure 3	9
Figure 4	10
Figure 5(a).....	11
Figure 5(b).....	12
Figure 5(c).....	13
Figure 6	14
Figure 7	17
Figure 8	21
Figure 9	22

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I. ABSTRACT

The non-linear inverse method of Xie (1993) is applied to analyze Lg spectra from 21 recent underground nuclear explosions and 52 shallow (5-33 km) earthquakes in central Eurasia. The data set used in this study consists hundreds of high-quality Lg spectra, collected from broad-band IRIS, CDSN and KNET stations. For those events from which Lg spectra at multiple (≥ 3) stations are available, the analyses simultaneously determine the Lg seismic moments (M_0), corner frequencies (f_c), path-variable Lg Q values at 1 Hz and their frequency dependences. For the events with Lg spectra recorded at only one or two stations, the analyses typically only determine Lg M_0 and f_c values, with path Lg Q values fixed using *a priori* information obtained from earlier inversions. The main findings of this study include

- (1) Lg Q values at 1 Hz for numerous paths in central Eurasia generally agree well with those predicted using a tomographic Lg coda Q map (Xie & Mitchell, 1991). This suggests that the 1 Hz Lg Q values obtained in this inversion have not been significantly biased by effects of non-isotropic source radiation patterns or large-scale 3D structural complications. The power-law frequency dependences of Lg Q and Lg coda Q also agree at distances between about 800 and 2700 km. At larger distances the Lg Q tends to show low (~ 0.0) frequency dependence.
- (2) For both underground nuclear explosions and the earthquakes studied, the logarithm of Lg M_0 values correlate linearly with the ISC body wave magnitude (M_b), with slopes of slightly greater than 1.0. For the same M_0 values, M_b values for the earthquakes tend to be systematically lower than for the explosions.
- (3) For the 21 underground nuclear explosions, the Lg M_0 values estimated using the explosion source model (*i.e.*, the model with an overshoot effect) scale with f_c^{-4} , instead with f_c^{-3} , as predicted by a constant stress drop scaling relationship.
- (4) For the earthquakes, the Lg M_0 values estimated using the earthquake source model (*i.e.*, the ω^2 model without overshoot) scale with $f_c^{-\alpha}$, with the value of α being about 3.6 when all of the 53 earthquakes are used in a regression analysis, and about

4.0 when only the well-recorded earthquakes (*i.e.*, the earthquakes with Lg spectra recorded at three or more stations) are used.

- (5) To simulate a situation in which we do not know that the explosions are explosions, we also inverted for their Lg M_0 and f_c values using the *earthquake* source model. The Lg M_0 and f_c values thus estimated are systematically different from those estimated for the explosions using the *explosion* source model: for the same f_c values, the Lg M_0 values estimated using the explosion source model is systematically lower (by a factor of 0.27) than those estimated using the earthquake source model. This factor indicates that there is a strong model-dependence in the estimated Lg M_0 and f_c values when the Lg from explosions is studied.
- (6) The scaling between Lg M_0 and f_c values for the explosions, estimated using the *earthquake* source model, also differ from that between the Lg M_0 and f_c values for the earthquakes. The main difference between the two scalings is that for a given M_0 , the explosions tend to have higher f_c values. This suggests that the Lg from explosions has a richer high-frequency content, as compared to Lg from earthquakes. This phenomenon appears to be systematic in central Eurasia, and appear to be opposite to that observed in the western U.S. It is therefore highly recommended that more research be conducted on which of these phenomena is more common for various continental environments.

II. STATEMENT OF WORK

Xie (1993) developed a non-linear inverse method for simultaneously estimating Lg source spectral parameters and path-variable Lg Q values. The advantages of that method, including its requiring no starting model for the unknown parameters and its allowing Lg Q to be path-variable, are fully described in Xie (1993). The proposed work for this two-year research is to apply the method of Xie (1993) to broad-band Lg records from many underground nuclear explosions and earthquakes occurring in central Asia after late 1987, the time when installations of broad-band stations began in Eurasia. The proposed steps of this study are:

- (1) Determine path variable Q_0 and η (Lg Q at 1 hz and its power-law frequency dependence; defined via $Q(f) = Q_0 f^\eta$) values from the Balapan and Lop Nor test sites to the recently installed broad-band seismic stations in Eurasia.
- (2) Compare the resulting Q_0 and η values with those predicted using the tomographic Lg coda Q map (Xie & Mitchell, 1991; Pan *et al.*, 1992). Find if these values disagree due to effects of either abnormal attenuation of Lg or Lg coda caused by large-scale 3D structural complications, or non-isotropic radiation patterns by the seismic sources.
- (3) Determine source spectral parameters, including seismic moment (M_0) and corner frequency (f_c), for explosions using the modified Mueller-Murphy source model, which is characterized by an ω^{-2} high-frequency asymptote and an overshoot effect:

$$S^{\text{exp}}(f) = \frac{M_0}{4\pi\rho v_s^3} \frac{1}{\left[1 + (1 - 2\beta)f^2/f_c^2 + \beta^2 f^4/f_c^4\right]^{1/2}} \quad (1)$$

(Serenio *et al.*, 1988; Xie, 1993) . Determine the scaling between the resulting M_0 and f_c values.

- (4) Determine M_0 and f_c values for earthquakes using the ω^{-2} earthquake source model with no overshoot effect:

$$S^{\text{eq}}(f) = \frac{M_0}{4\pi\rho v_s^3} \frac{1}{1 + f^2/f_c^2} \quad (2)$$

(Street *et al.*, 1975; Xie, 1993). Determine the scaling between the resulting M_0 and f_c values.

- (5) From the results of (3) and (4), determine if M_0 scales with f_c^{-3} or f_c^{-4} for explosions and earthquakes, and infer the resulting Lg source spectral scaling in terms of any fundamental differences between the excitation of Lg and that of local S waves by

both types of seismic sources (*i.e.*, to infer if the transfer function between the Lg and S excitations is flat; for details, *cf.* later sections).

- (6) Determine the M_0 and f_c values for explosions using the *earthquake* source model, thus simulating a situation when we do not know that the events are explosions. Determine the scaling between the resulting M_0 and f_c values.
- (7) Compare the the scaling derived in step (6) with that obtained in step (4) to see the difference between the $M_0 - f_c$ scalings for explosions and earthquakes, both derived using the earthquake source model. Determine if the difference can be used as an effective discriminant of explosions.

III. RESEARCH ACCOMPLISHED

Over the past two years we have collected hundreds of Lg spectra from 21 underground nuclear explosions and 53 shallow (5-33 km) earthquakes in/around the Balapan and Lop Nor test sites, recorded by 21 broad-band IRIS, CDSN and KNET stations. All events studied occurred after 1987, the time when the installation of broad-band digital stations began in Eurasia. Tables 1 and 2 lists the explosions and earthquakes studied, respectively. Figure 1 shows the locations of the events and stations used in this study. The earthquakes are chosen to be either near the test sites, or near the KNET stations, so that the source locations or paths involved in studying the earthquakes are similar to those involved in studying the explosions. The lengths of the paths used in this study are between about 800 km and 4045 km (see Figures 5a through 6).

Data processing

The Lg spectra are obtained in the same way as that described by Xie (1993). Figure 2 shows an example of the Lg time series, the 20% taper window used, and the resulting Lg spectra. For each event studied, we used the non-linear inverse method by Xie (1993) to invert for source and path spectral parameters. When the event is recorded by a sufficient number of stations (≥ 3), we simultaneously inverted for source M_0 , f_c and path Q_0 , η values with no *a priori* information. For the events that are recorded by only one or two stations, we found that the information contained in the Lg spectra is typically not sufficient to constrain a simultaneous inversion and accordingly, for most of these events we used *a priori* information on path Q_0 and η values obtained in the earlier simultaneous inversions. Figures 3 and 4 show examples of the fit of the optimal source/path parameters to the observed Lg spectra, where the inverted Q_0 , η values for multiple stations that recorded an explosion (Figure 3) and earthquake (Figure 4) are used to remove path

Table 1. Underground Nuclear Explosions Studied†

Event Date	Origin Time	m_b	Test Site	Seismic Moment (Nm)	Corner Frequency	Number of Stations Used
Dec. 27, 87	03H05M04.9S	6.1	Balapan	$1.3 (\pm 0.3) \times 10^{16}$	0.61 ± 0.04 Hz	2
Feb. 13, 88	03H05M05.9S	6.1	Balapan	$1.7 (\pm 0.6) \times 10^{16}$	0.62 ± 0.05 Hz	1
Apr. 3, 88	01H33M05.8S	6.0	Balapan	$1.7 (\pm 0.3) \times 10^{16}$	0.59 ± 0.04 Hz	2
May 4, 88	00H57M06.8S	6.1	Balapan	$1.9 (\pm 0.6) \times 10^{16}$	0.59 ± 0.05 Hz	1
Jun. 14, 88	02H27M06.4S	5.1	Balapan	$7.0 (\pm 3.4) \times 10^{14}$	1.20 ± 0.11 Hz	1
Sep. 14, 88*	03H59M57.6S	6.1	Balapan	$1.3 (\pm 0.1) \times 10^{16}$	0.56 ± 0.02 Hz	5
Nov. 12, 88	03H30M30.8S	5.4	Balapan	$2.6 (\pm 0.3) \times 10^{15}$	0.82 ± 0.02 Hz	3
Nov. 23, 88	03H57M06.8S	5.4	Balapan	$1.8 (\pm 0.2) \times 10^{15}$	0.70 ± 0.02 Hz	3
Dec. 17, 88	04H18M06.8S	5.9	Balapan	$1.3 (\pm 0.4) \times 10^{16}$	0.52 ± 0.03 Hz	4
Jan. 22, 89	03H57M06.7S	6.0	Balapan	$1.2 (\pm 0.3) \times 10^{16}$	0.60 ± 0.04 Hz	3
Feb. 12, 89	04H15M06.9S	5.8	Balapan	$9.5 (\pm 2.5) \times 10^{15}$	0.59 ± 0.03 Hz	4
Jul. 8, 89	03H47M00.0S	5.6	Balapan	$3.5 (\pm 1.2) \times 10^{15}$	0.70 ± 0.05 Hz	2
Sep. 2, 89	04H16M59.9S	5.1	Balapan	$8.2 (\pm 3.4) \times 10^{14}$	1.01 ± 0.07 Hz	1
Oct. 19, 89	09H49M58.0S	5.9	Balapan	$1.0 (\pm 0.2) \times 10^{16}$	0.64 ± 0.04 Hz	5
Aug. 16, 90	04H59M57.7S	6.2	Lop Nor	$7.9 (\pm 1.4) \times 10^{14}$	0.54 ± 0.02 Hz	2
May 21, 92	04H59M57.5S	6.6	Lop Nor	$4.3 (\pm 0.7) \times 10^{16}$	0.40 ± 0.02 Hz	4
Oct. 5, 93	03H59M57.6S	5.9	Lop Nor	$8.3 (\pm 0.2) \times 10^{15}$	0.68 ± 0.05 Hz	7
Jun. 10, 94	06H25M58.0S‡	5.7‡	Lop Nor	$2.0 (\pm 0.4) \times 10^{15}$	0.92 ± 0.03 Hz	5
Oct. 7, 94	03H25M58.0S‡	5.9‡	Lop Nor	$5.5 (\pm 0.2) \times 10^{15}$	0.80 ± 0.06 Hz	8
May 5, 95	04H05M58.0S‡	5.9‡	Lop Nor	$7.8 (\pm 0.2) \times 10^{15}$	0.66 ± 0.04 Hz	8

† The origin times, locations and magnitudes are from the ISC or PDE bulletin; the seismic moments and corner frequencies are obtained in this study using the explosion source model with β (equation (2)) set to be 0.75 (Poisson medium).

* The Joint Verification Experiment (JVE) event.

‡ USGS preliminary (PDE) estimate.

Table 2. Earthquakes Studied†

Number	Date	Origin Time h m s	Latitude (° N)	Longitude (° E)	m_b	M_0 (dyne-cm)	f_c (Hz)	Number of Stations
1	03 Jan 88	20 09 21.4	38.431	91.340	4.4	1.9×10^{22}	0.84	1
2	06 Feb 88	04 19 11.1	49.799	78.064	4.8	7.0×10^{21}	0.65	1
3	15 Mar 88	15 55 24.3	42.210	75.509	4.5	2.2×10^{22}	0.50	1
4	05 Mar 89	13 48 41.6	42.511	74.629	4.8	7.0×10^{22}	0.60	1
5	14 Apr 89	22 57 59.6	41.132	74.525	4.6	1.4×10^{22}	0.50	1
6	21 Jan 90	07 53 31.9	41.534	88.728	4.6	1.0×10^{22}	0.80	1
7*	02 Feb 90	14 04 25.5	42.219	76.270	4.4	1.4×10^{22}	0.76	2
8*	03 May 90	10 02 22.2	42.790	76.880	4.7	3.8×10^{22}	0.70	3
9	19 Sep 90	08 05 57.3	38.001	88.940	4.4	1.8×10^{22}	0.90	1
10*	03 Nov 90	17 25 13.8	40.882	89.071	5.1	1.2×10^{23}	0.50	3
11	08 Jun 92	09 20 54.5	43.598	88.277	4.2	5.0×10^{21}	1.00	1
12	10 Jun 92	02 37 01.2	38.623	90.147	4.4	1.6×10^{22}	0.88	1
13*	19 Aug 92	10 17 35.2	42.265	73.252	5.1	2.2×10^{23}	0.38	4
14	19 Aug 92	14 17 40.7	41.876	73.410	4.7	6.1×10^{22}	0.42	1
15*	19 Aug 92	22 45 51.2	41.897	73.199	4.9	6.0×10^{22}	0.60	2
16	20 Aug 92	01 28 02.5	41.751	73.361	4.6	2.8×10^{22}	0.66	1
17	20 Aug 92	06 25 47.0	41.864	73.386	4.5	3.2×10^{22}	0.63	1
18	20 Aug 92	06 52 41.8	41.951	73.215	4.1	1.8×10^{22}	0.75	1
19	20 Aug 92	10 01 16.5	41.699	73.086	4.0	1.8×10^{22}	0.70	1
20	20 Aug 92	12 22 47.3	41.991	73.377	4.8	1.0×10^{23}	0.57	2
21	20 Aug 92	12 59 27.7	41.836	73.058	4.2	1.9×10^{22}	0.88	1
22	20 Aug 92	16 30 45.4	41.837	73.633	4.2	2.5×10^{22}	0.88	1
23	20 Aug 92	21 35 30.0	42.212	73.507	4.5	1.2×10^{22}	0.87	1
24	21 Aug 92	04 14 32.7	41.923	73.503	4.7	7.1×10^{22}	0.66	2
25	22 Aug 92	08 52 08.2	41.996	73.470	4.4	2.3×10^{22}	0.76	1
26	23 Aug 92	00 28 06.2	41.950	73.746	4.1	1.1×10^{22}	0.88	1
27	23 Aug 92	02 54 47.4	41.892	73.575	4.4	1.7×10^{22}	0.77	1
28	23 Aug 92	07 15 07.3	41.909	73.463	4.7	4.0×10^{22}	0.62	1
29	23 Aug 92	09 04 32.4	41.998	73.567	4.9	8.5×10^{22}	0.70	2
30	23 Aug 92	20 11 42.4	41.951	73.447	4.4	2.1×10^{22}	0.75	1
31	23 Aug 92	20 35 06.0	41.959	73.536	4.6	5.1×10^{22}	0.61	1
32	26 Aug 92	07 40 36.7	41.785	73.387	4.8	4.2×10^{22}	0.73	1
33	26 Aug 92	20 44 40.3	41.941	73.664	4.4	1.0×10^{22}	0.66	1
34	26 Aug 92	22 01 15.0	41.928	73.552	4.6	3.0×10^{22}	0.70	1
35	28 Aug 92	04 33 38.8	41.933	74.395	4.5	1.1×10^{22}	0.69	1
36	25 Sep 92	07 59 59.9	41.763	88.387	5.0	1.2×10^{22}	0.82	1
37	20 Oct 92	16 30 52.0	41.927	73.241	4.0	7.0×10^{21}	1.08	1
38*	27 Nov 92	16 09 09.1	41.978	89.283	5.3	1.4×10^{23}	0.48	3
39*	02 Feb 93	16 05 14.1	42.219	86.132	5.7	3.7×10^{23}	0.34	4
40	17 Feb 93	02 00 25.8	38.321	89.484	5.1	1.5×10^{23}	0.58	1
41	13 Apr 93	17 56 02.0	41.190	75.719	4.7	1.1×10^{23}	0.46	1
42	14 Apr 93	08 31 09.7	42.904	87.045	4.4	1.3×10^{22}	0.64	1
43	26 May 93	14 11 12.4	40.117	91.525	4.4	8.0×10^{21}	0.82	2
44*	02 Oct 93	08 42 32.7	38.190	88.663	6.2	4.8×10^{24}	0.21	7
45	02 Oct 93	09 20 12.2	38.206	89.284	4.9	4.7×10^{22}	0.62	1
46*	02 Oct 93	09 43 19.5	38.169	88.605	5.8	4.8×10^{23}	0.33	5
47*	02 Oct 93	17 23 33.3	38.171	88.690	5.6	2.6×10^{23}	0.41	6
48	02 Oct 93	19 16 43.0	38.079	88.831	3.8	7.4×10^{21}	1.10	1
49	02 Oct 93	23 49 59.7	38.359	88.878	4.8	1.6×10^{22}	0.64	1
50	07 Oct 93	03 26 58.9	38.214	88.726	5.0	3.4×10^{22}	0.68	2
51	12 Oct 93	20 49 23.4	38.276	88.604	4.7	5.6×10^{22}	0.54	1
52	08 Jun 94	21 03 41.4	43.228	86.886	4.6	1.2×10^{22}	0.69	1

† The origin times, locations and magnitudes are from the IRIS DMC; the seismic moments (M_0) and corner frequencies (f_c) are obtained in this study using the earthquake source model.

* Events for which full inversion for M_0 , f_c , Q and η for paths to station was performed. (for other events, Q was fixed at values determined along similar paths from earlier inversion).

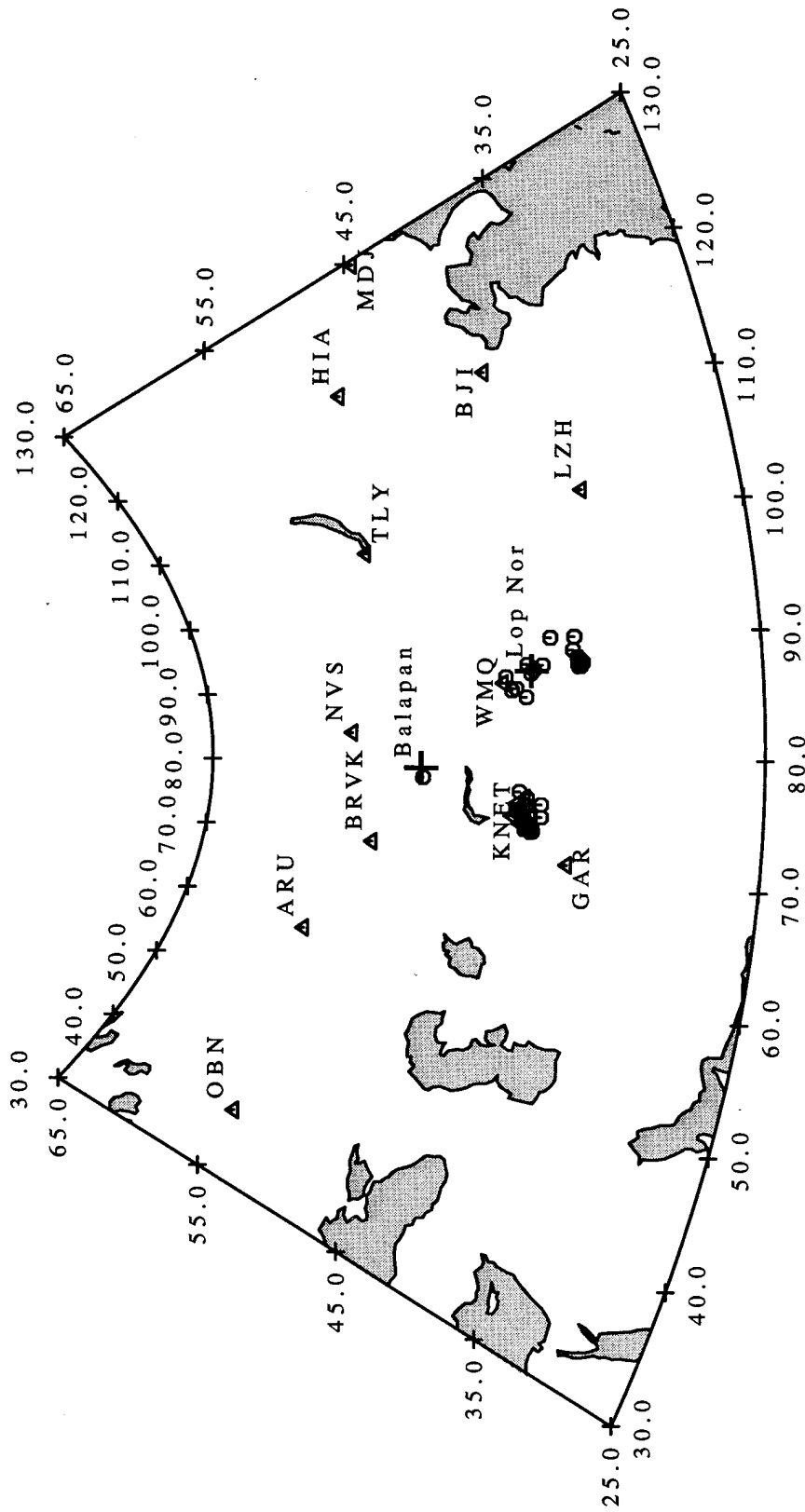


Fig. 1. Locations of the 21 underground nuclear explosions (crosses), 52 earthquakes (circles) and 21 seismic stations (triangles) used in this study. Water covered areas are shaded. The numbers of stations providing Lg records are 17 for the explosions, and 11 for the earthquakes.

Station: CHM. Starting time=282.4sec. Dist=1157.8km

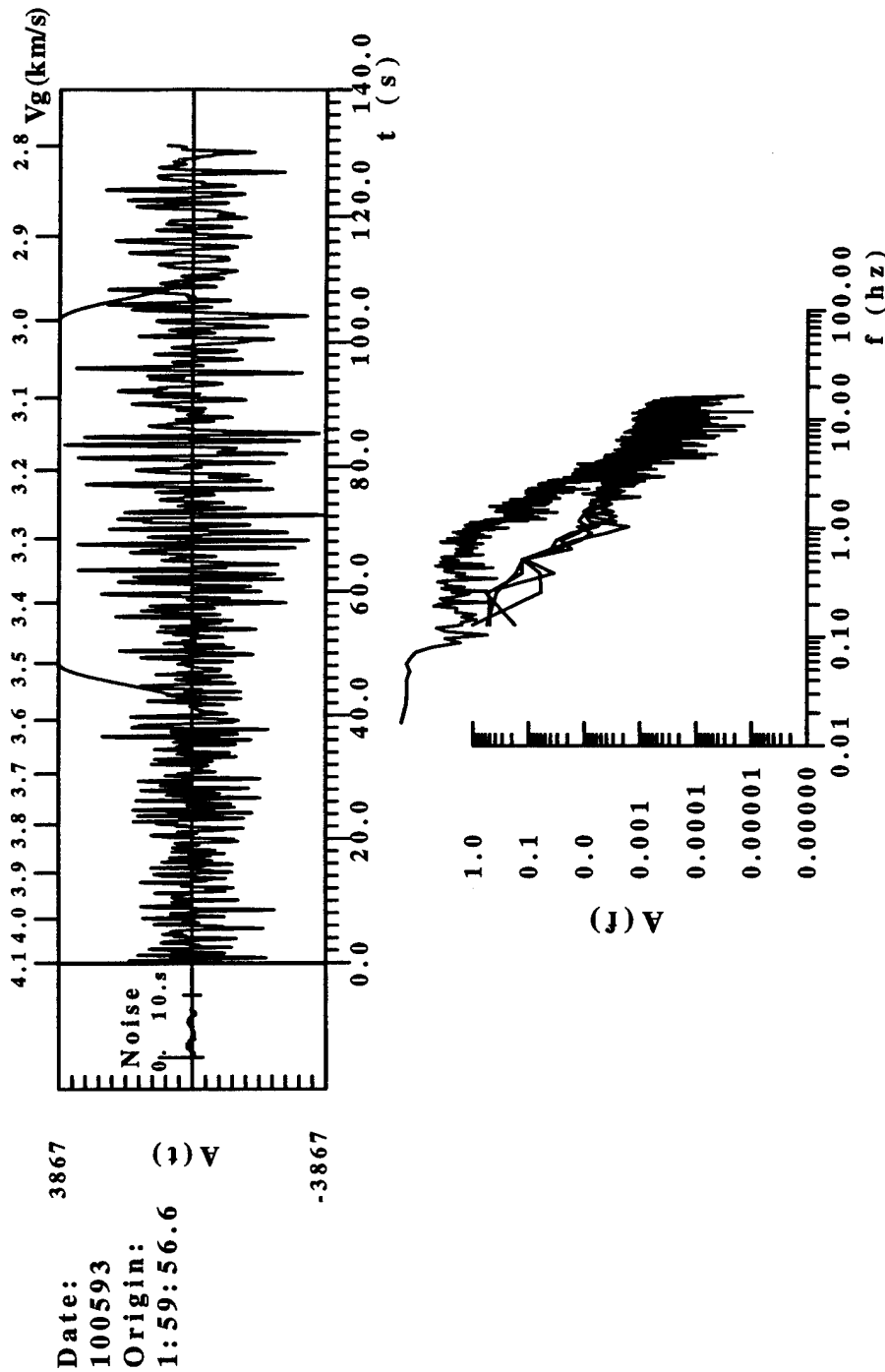


Fig. 2. (Top) Time series containing Lg from a Lop Nor explosion (see left of the panel) at KNET station CHM. Group velocities and time after the beginning of the window are indicated on the top and bottom of the panel. The smooth curve represents a 20 percent cosine window used in the analysis. (Bottom) The instrument-corrected Fourier amplitude spectra of Lg and noise prior to P.

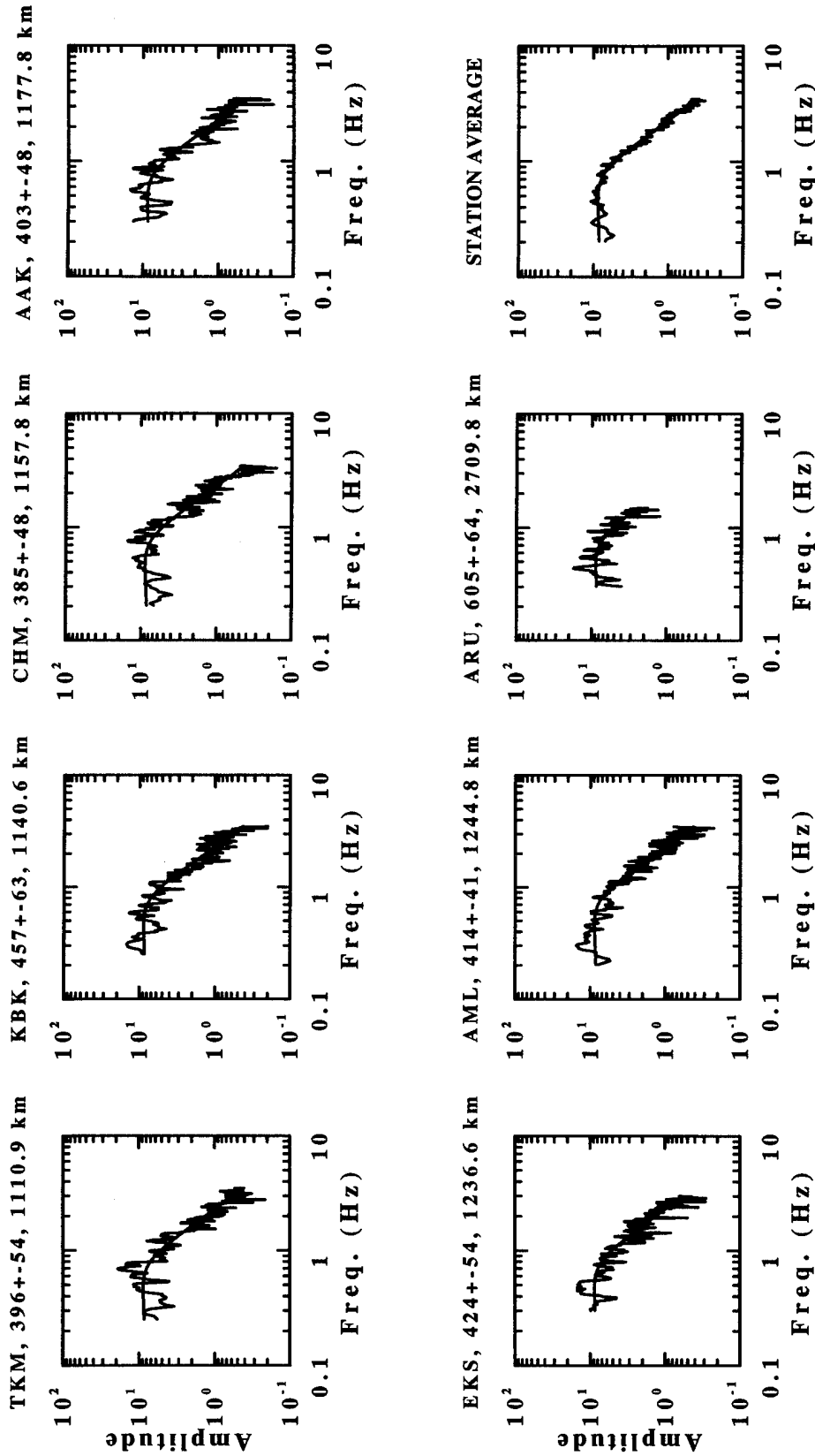


Fig. 3. Synthetic Lg source spectra for seven KNET and IRIS stations recording the October 5, 1993, Lop Nor explosion, versus the observed Lg spectra that are reduced to source by removing path effects. The lower right panel is the average for all of the seven stations. The synthetic spectra are calculated using optimal source spectral parameters ($M_0 = 8.3 \times 10^{15}$ Nm, $f_c = 0.68$ Hz) obtained in the inversion. Path Q_0 values obtained in the inversion are written on the top of the panels, together with the epicentral distances.

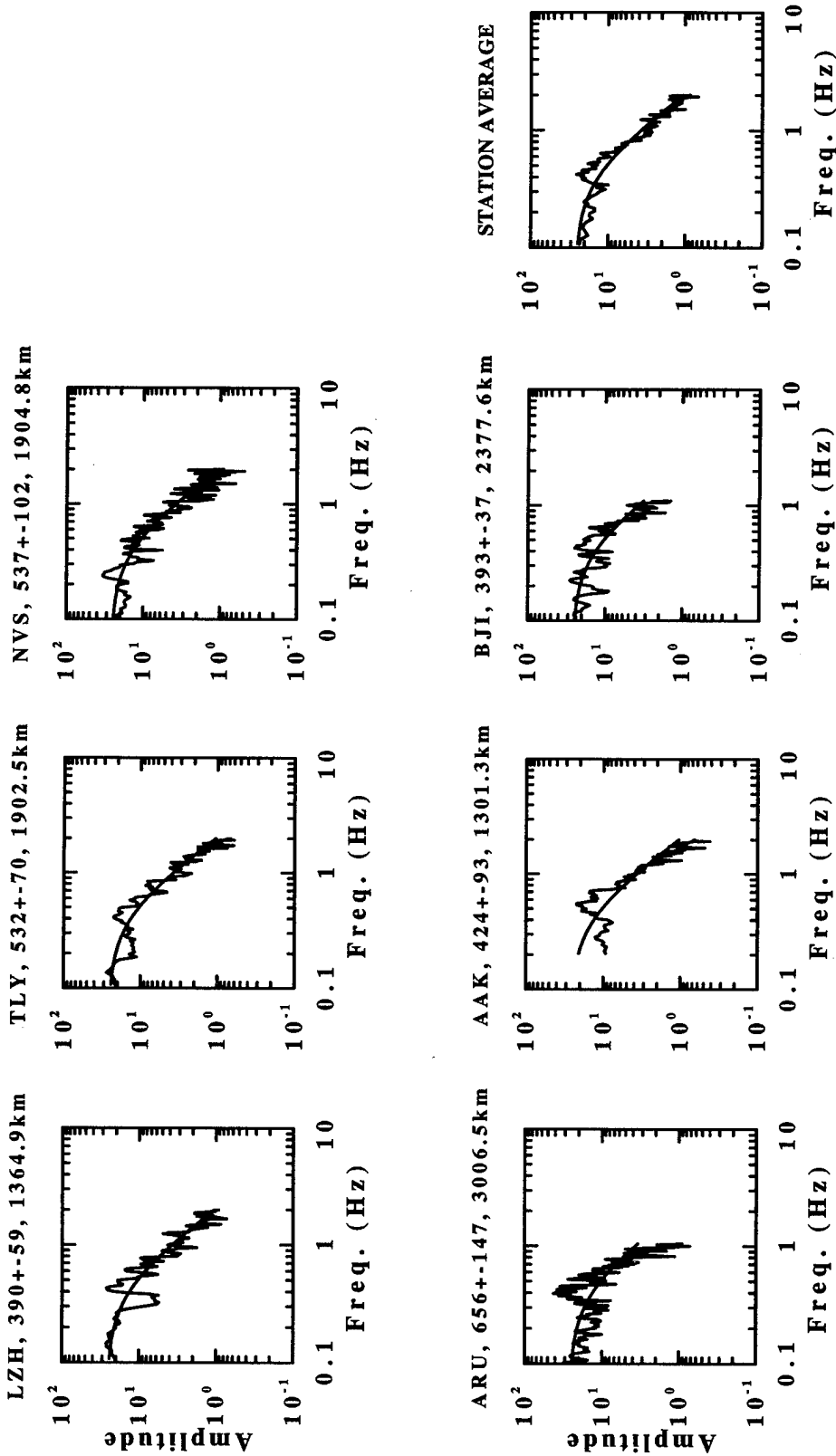


Fig. 4. Synthetic Lg source spectra for six IRIS and CDSN stations recording the October 2, 1993, southern Xinjiang earthquake ($M_b = 5.6$), versus the observed. The synthetic spectra are calculated using optimal source spectral parameters ($M_0 = 2.6 \times 10^{16}$ Nm, $f_c = 0.41$ Hz) obtained in the inversion.

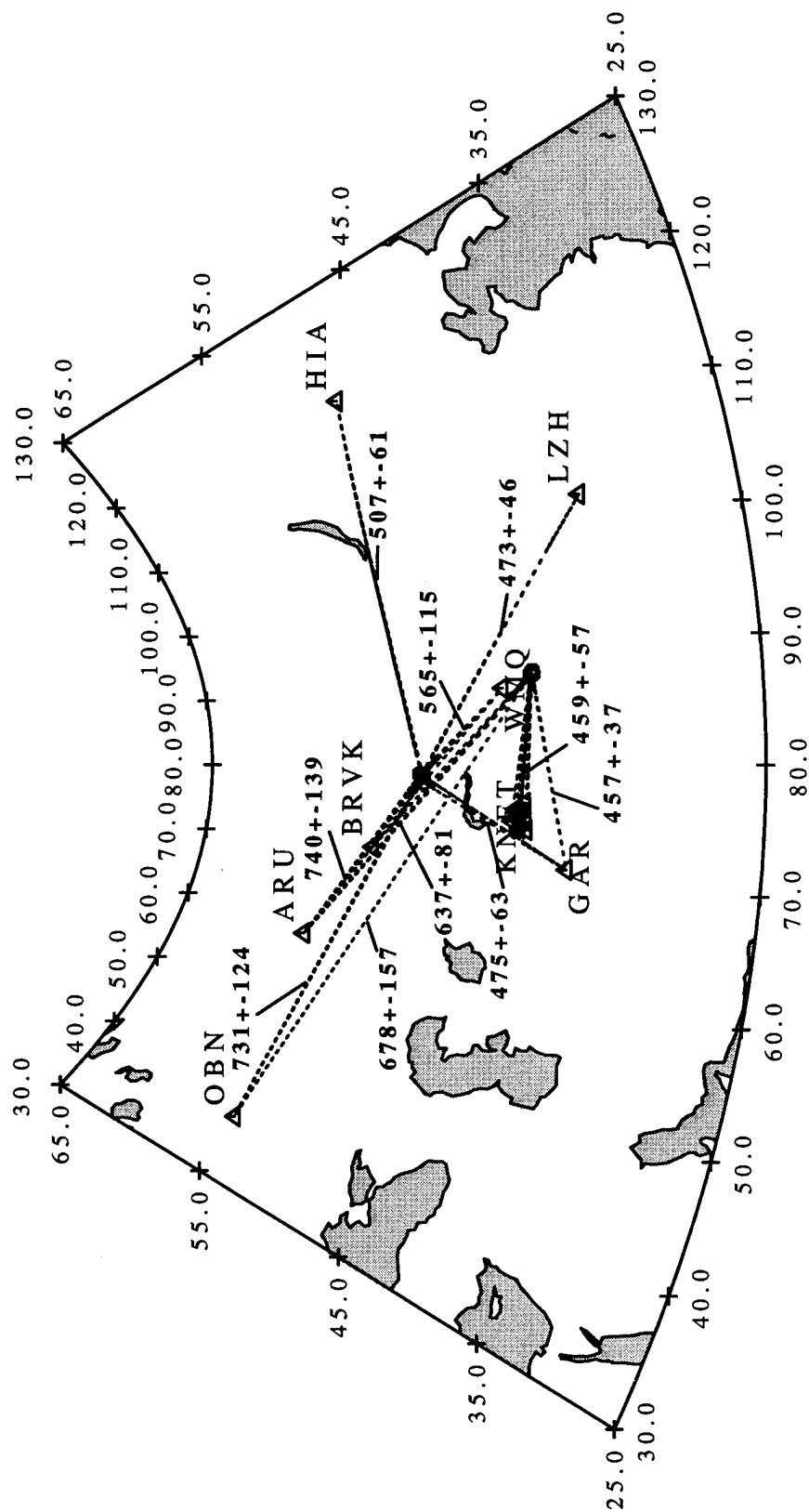


Fig. 5 (a) $Lg Q_0$ values obtained for the great circle paths from the Lop Nor and Balapan test sites to the 17 IRIS, CDSN and KNET stations. Water-covered areas are shaded.

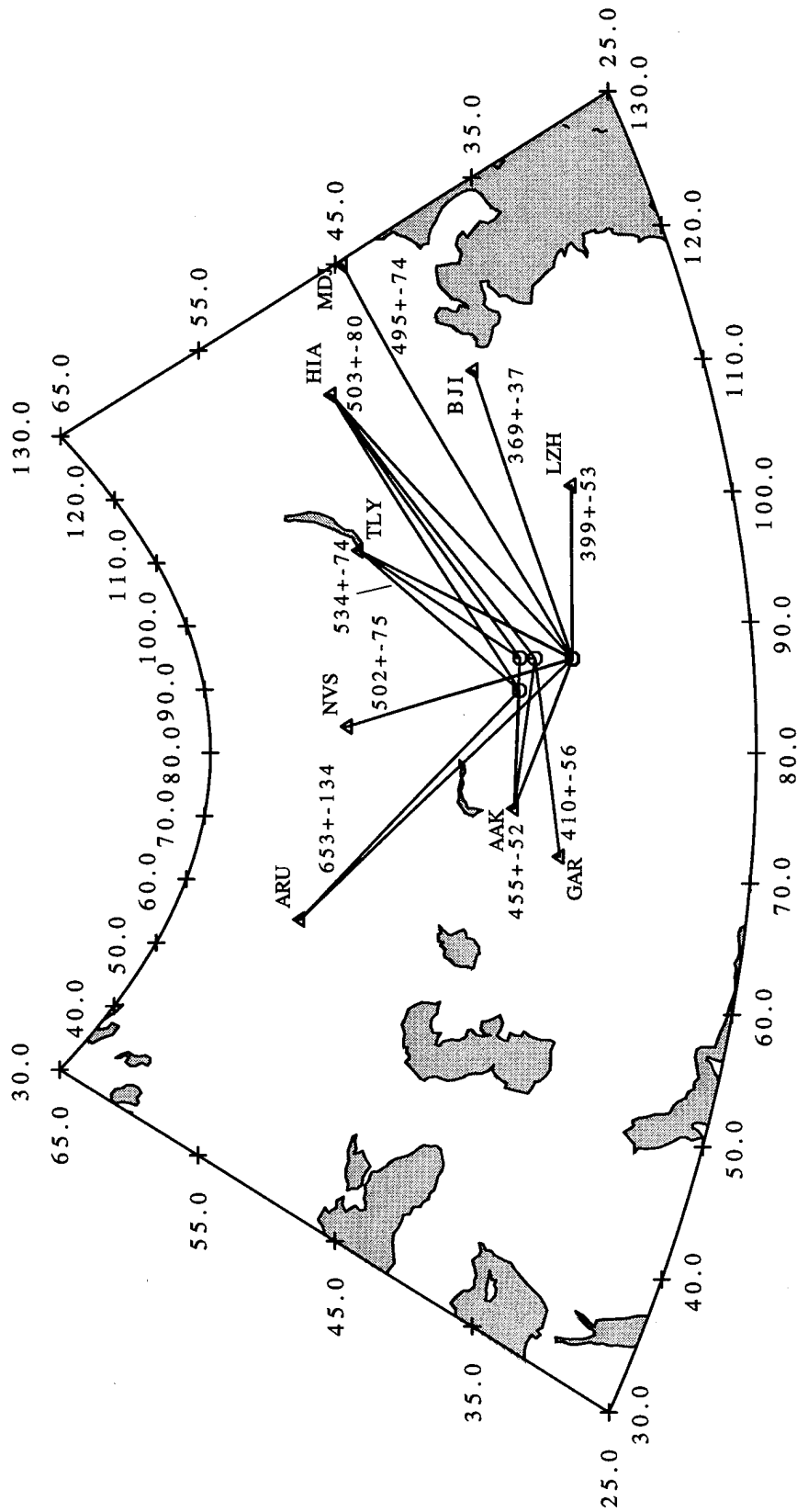


Fig. 5(b) Lg Q_0 values obtained for the great circle paths from earthquakes in Xin Jiang, China to IRIS and CDSN stations.

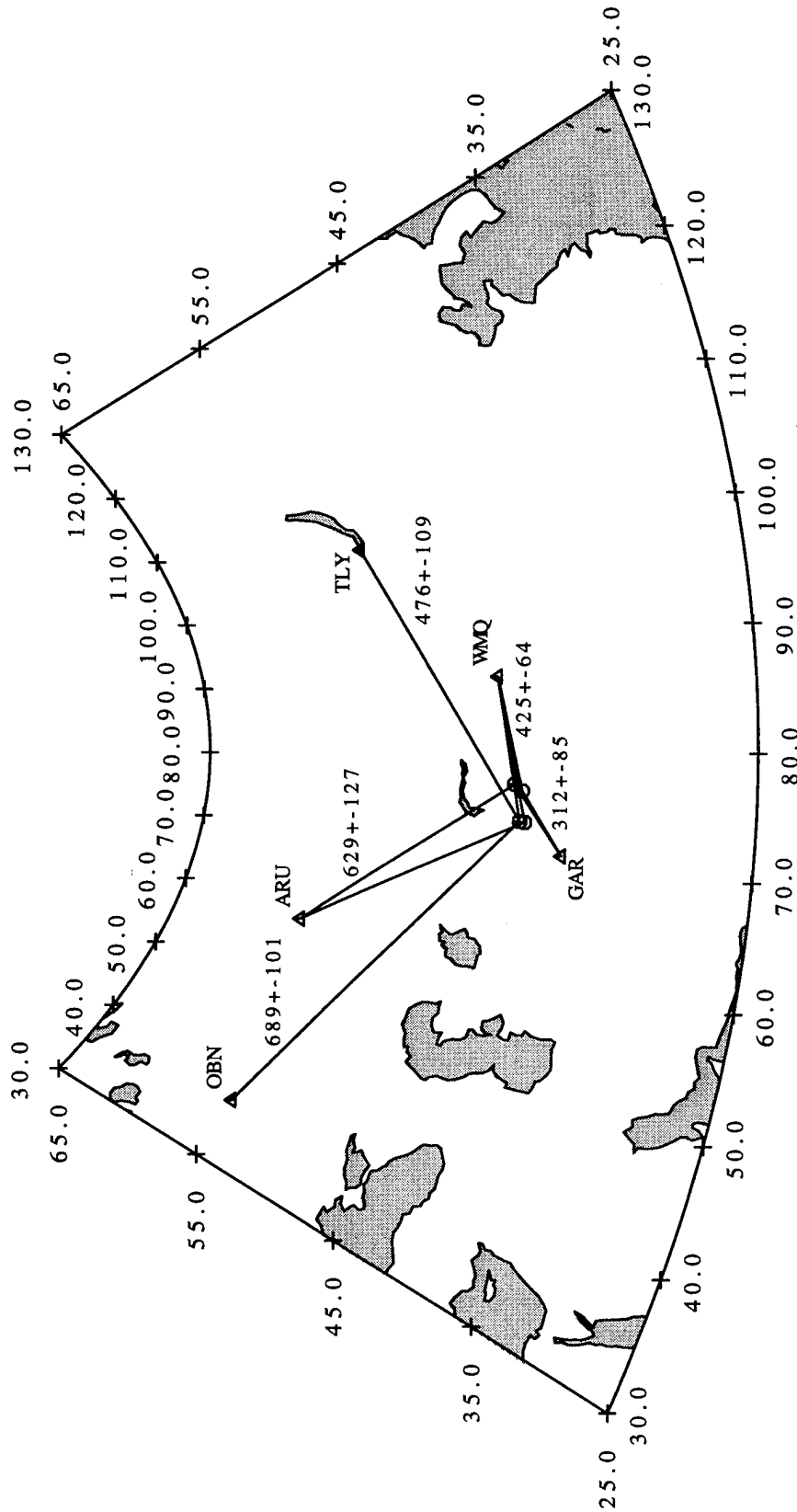


Fig. 5(c) Lg Q_0 values obtained for the great circle paths from earthquakes in the central Asian Republics of the F.S.U. to IRIS and CDSN stations.

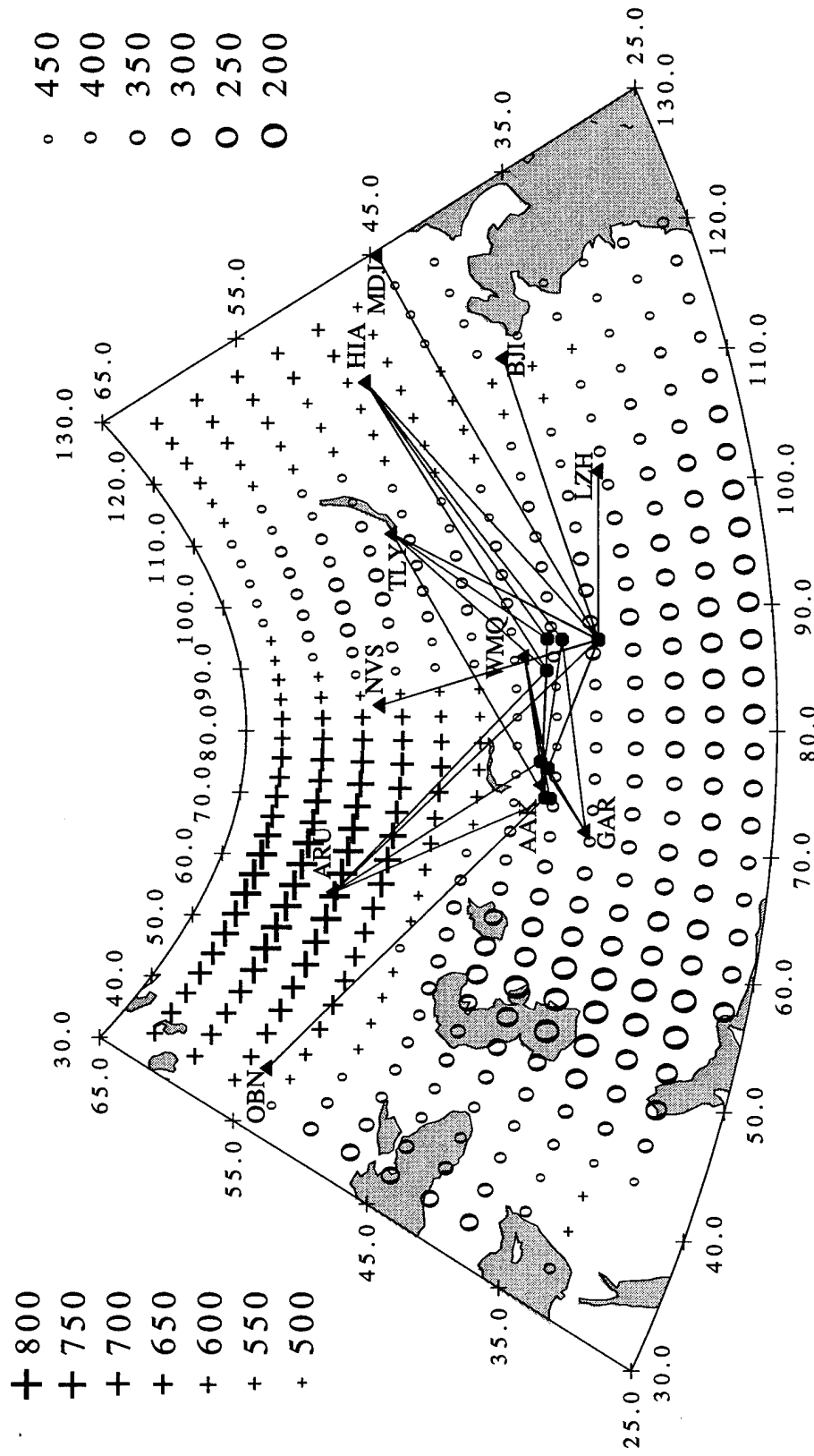


Fig. 6. The Lg paths for which Q values are measured in this study (see Figures 5a through 5c), plotted with the Lg coda Q_0 values from the tomographic inversion by Xie and Mitchell (1991). Solid dots and triangles are sources and stations, respectively.

effects, resulting in reduced Lg spectra at the source. The fit between the observed Lg spectra (reduced to source) and the theoretical source spectrum for the explosion event (Figure 3) and that for the earthquake event (Figure 4) are both good, particularly for the station averages (lower right panels). The fit in Figure 3, though, appear to be better than that in Figure 4. This phenomenon is typical for all events studied, indicating a greater complexity in Lg source spectra of earthquakes than those of explosions in the study area.

Agreements among 1 Hz Q measured using Lg from explosions, Lg from earthquakes and Lg coda

In several previous studies (*eg.*, Kopnichev, 1977; Herrmann, 1980; Der *et al.*, 1984; Xie & Nuttli, 1988; Xie & Mitchell, 1990b; Ryaboy, 1990; Chun *et al.*, 1994), observed Lg coda Q values were similar to those of Lg Q when both were carefully measured. Considering the 3D complexity of the earth, however, the agreement between Lg Q and Lg coda Q is somewhat surprising. Regions where they do not agree are usually sites of major lateral crustal discontinuities; thus discrepancy between Lg Q and Lg coda Q maybe a tool to detect large-scale disruptions of the crustal waveguide (Xie & Mitchell, 1990b, Kennett *et al.*, 1991). In this study we have obtained Lg Q_0 and η values for numerous paths covering a large area of central Eurasia. In Figures 5a through 5c we plotted the Lg Q_0 values (Lg Q at 1 Hz) obtained for all of the path groups obtained in this study. To compare these Lg Q_0 values with Lg coda Q_0 values, in Figure 6 we plotted a portion of the tomographic Lg coda Q_0 map by Xie & Mitchell (1991). The estimated errors of Q_0 values in Fig. 6 are about 10% to 15%. Similarities in the Q_0 values for each path in Figure 5a and those predicted for the same path in Fig. 6 are readily apparent. For example, from Lop Nor to the KNET stations, Fig. 5a gives an average Q_0 of 459 ± 57 whereas from Fig. 6 we can predict a Q_0 around 450. From Lop Nor to stations ARU and OBN, Fig. 5a gives Q_0 of 637 ± 81 and 678 ± 157 , respectively. These agree with corresponding predictions of around 650 and 700 based on Fig. 6. Agreement can also be found for paths from Balapan to various IRIS/CDSN stations. Xie (1993) has already discussed the latter agreement, using results from analyzing a much smaller data base. Since the explosion sources should have near-isotropic moment tensors, it is reasonable to assume that the source to Lg radiation does not vary significantly with azimuth. The agreement between the Q_0 values in Figures 5a and 6 indicates, then, that the Q_0 values obtained in this study using Lg from explosions have not been seriously biased by effects of large-scale 3D structural complications, such as the focusing/defocusing effects by the Tianshan mountain predicted by Bostock & Kennett (1990).

For earthquake sources, the possibility of non-isotropic source radiation pattern is expected to increase, and there is a concern that the Q measured using Lg from earthquakes may be biased by that radiation pattern. We note that the paths connecting the earthquakes near the Lop Nor test site and stations ARU, GAR and AAK in Figure 5b roughly overlap with the paths connecting the Lop Nor explosions and the respective stations in Figure 5a. For these three paths, Figure 5b shows Q_0 values of 653 ± 134 (to ARU), 410 ± 56 (to GAR), and 455 ± 52 (to AAK), respectively. These values agree, within the estimated uncertainties, with the corresponding values of 637 ± 81 , 457 ± 37 , and 459 ± 57 in Figure 5a. Therefore for all of the available overlapping paths, we find no difference between Q_0 measured using Lg from earthquakes and Q_0 measured using Lg from explosions.

For other paths in Figures 5b and 5c, no overlapping paths can be found in Figure 5a to warrant a comparison in Q_0 , but we can still compare the Lg Q_0 values for these paths with those predicted by the Lg coda Q_0 map in Figure 6. In general, the Q_0 values in Figures 5b and 5c are very compatible to those in Figure 6. For example, the Lg Q_0 value along the path to station LZH in Figure 5b is 399 ± 53 , which matches the values of about 350 to 400 in Figure 6. Similarities between Q_0 values for other paths in Figures 5b or 5c and those in Figure 6 are readily apparent, indicating that the Lg Q_0 measured using Lg from earthquakes are similar to those predicted by the Lg coda Q_0 map.

We conclude that all comparisons available from this study indicate that the Q_0 values measured using Lg from explosions agree with Q_0 values measured using Lg from earthquakes, and these Lg Q_0 values agree with previously mapped Lg coda Q_0 values in the area. Although these agreements do not preclude possibilities of observing any future disagreement between Lg Q_0 and Lg coda Q_0 , it does indicate that the Lg Q_0 values obtained in this study have not been seriously biased by effects of any non-isotropic source radiation patterns, or of any large-scale 3D structural complications, such as the focusing/defocusing effects by the Tianshan mountain predicted by Bostock & Kennett (1990).

Comparison between Lg η and Lg coda η

We also compared the values of Lg η , the power-law frequency dependence of Lg Q , with the values predicted with an Lg coda η map (Xie and Mitchell, 1991). We found that when the epicentral distance (Δ) is less than about 2700 km, the Lg η obtained in this study agrees (within an uncertainty level of about 0.1 to 0.2) with the Lg coda η . At larger distances ($\Delta > 2700$ km), the η values obtained in this study tend to be low (typically down to ~ 0.0). This discrepancy may be caused by a number of reasons, including

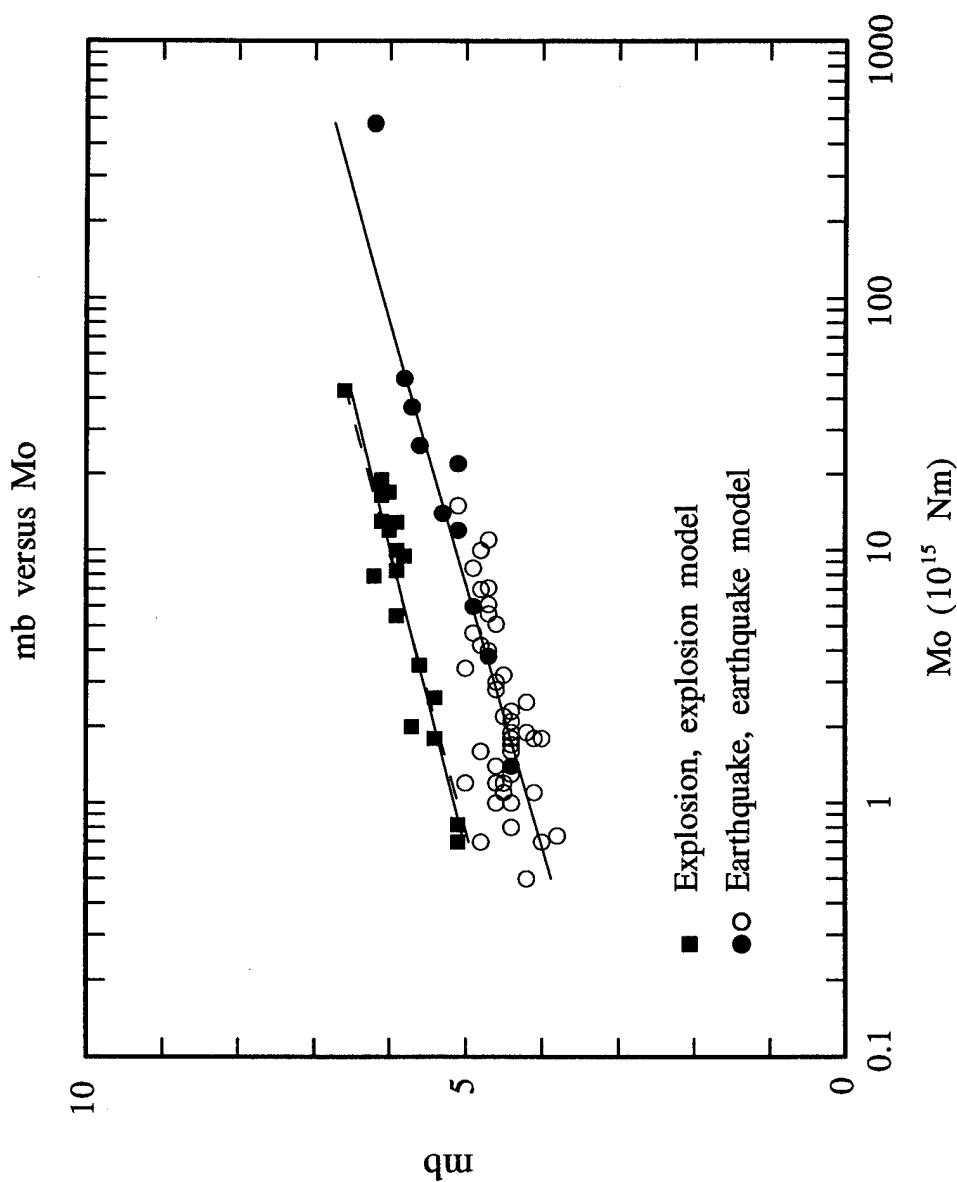


Fig. 7. M_b values versus logarithm of M_0 values (in 10^{15} Nm) obtained for explosions (squares) and earthquakes (circles). Solid circles represent values for earthquakes that are better recorded (*i.e.*, earthquakes with Lg recorded at three or more stations), and open circles represent values for earthquakes with Lg recorded at only one or two stations. Straight lines represent the linear regression fitting.

imprecisely estimated Lg η due to the narrower frequency bands, or effects of the earth's curvature at large Δ . Other possible reasons for the discrepancy is being explored but can not be currently identified with confidence.

Correlation between $M_0(\text{Lg})$ and ISC M_b

Figure 7 shows the M_0 values obtained for the explosions (squares) and earthquakes (circles). The M_0 values for the explosions are obtained with the explosion source model (equation (1)), and are the first ever obtained using Lg from underground nuclear explosions. We would therefore like to assess their reliability and consistency with the P wave seismic moments. For P waves the seismic moment, M_0^P , is related in theory to the displacement potential, ψ_∞ , via

$$\psi_\infty = \frac{M_0^P}{4\pi\rho\alpha^2} \quad , \quad (3)$$

where ρ and α are the source-zone density and P wave velocity (Mueller, 1973; Aki *et al.*, 1974). Ringdal *et al.* (1992) obtained an empirical relationship between ψ_∞ and ISC body-wave magnitude, M_b :

$$\log \psi_\infty = 1.1M_b + -2.57 (+ - 0.11) \quad (4)$$

(equation (13) of Ringdal *et al.*, 1992). Substitute equation (3) into (4) we have

$$\log M_0^P = 1.1M_b - 2.57 - \log(4\pi\alpha^2) (+ - 0.11) \quad . \quad (5)$$

Assuming that $\alpha = 5.2$ km/s and $\rho = 2.7$ g/cm³ (Li *et al.*, 1995), equation (3) becomes

$$\log M_0^P = 1.1M_b + 9.39 (+ - 0.11) \quad , \quad (6)$$

where M_0^P is in Nm. Equation (6) is empirical, and approximately relates the M_0^P to ISC M_b . To see if M_0^P predicted by equation (6) and the M_0 derived for the explosions using Lg in this study are consistent, we plotted equation (6) in Figure 7 (dashed line). The agreement between M_0^P values predicted by equation (6) (dashed line) and the M_0 values obtained in this study (in the M_b range between roughly 5.0 and 6.5) is very good. A linear regression over the Lg M_0 and m_b values for the explosions in Fig. 7 yields

$$\log M_0 = 1.19(+ - 0.11) M_b + 8.85(+ - 0.64) \quad (7)$$

for M_0 derived using Lg. The slope and intercept in (7) agrees with, within the uncertainties, those predicted by (6), even though the intercepts in both equations represent *extrapolations* of the linear trends to zero M_b .

Earlier observations have suggested that for explosions, the Lg magnitude (based on time domain amplitude measurement) correlated well with M_b (eg., Nuttli, 1986a, 1986b, 1988; Henson *et al.*, 1990; Ringdal *et al.*, 1992). These observations emphasize that both M_b and M_{bLg} are obtained by measuring short period (~ 1 Hz) amplitudes. The M_0 values in this study, on the other hand, are obtained by treating Lg as multiple supercritically reflected S waves (Street *et al.*, 1975; Xie, 1993). Furthermore, the M_0 values are mainly constrained by Lg spectra at lower frequencies (down to 0.1 to 0.2 Hz for the effective Lg pass-band). Thus the agreement between M_0^P and M_0 obtained in this study suggests a consistency between the source size measured using multi-station average of short-period P amplitudes and that using longer period multiple supercritically reflected S waves.

The Lg M_0 versus m_b for the 53 earthquakes in Figure 7 are subdivided into two groups: values represented by solid circles are for the earthquakes with Lg recorded at at least three stations, whereas values represented by open circles are for the earthquakes that are recorded by only one or two stations. For the latter group of earthquakes, we typically only inverted the M_0 and f_c values, with path Q_0 , η values fixed using *a priori* knowledge from other inversions. The solid line through the circles in Figure 7 represents a linear regression, which yields the following equation:

$$\log M_0 = 1.04(+ - 0.09) m_b + 10.66(+ - 0.52) \quad . \quad (8)$$

In Figure 7, the open circles show higher degree of deviation from the linear trend than shown by the solid circles; this is to be expected since the open circles correspond to events that are relatively poorly recorded.

When comparing the circles and squares in Figure 7, or equivalently, equations (7) and (8), the most important feature we find is that for the same M_0 values, the m_b values from explosions tend to be systematically larger than those from the earthquakes. Physically, this means that at the same Lg moment level, the explosions tend to have higher body wave amplitudes around 1 Hz than the earthquakes.

Previously unresolved issues on the scaling between M_0 and f_c values

The nature of the scaling between M_0 and f_c derived using the Lg source spectra have been controversial for quite a long time. For earthquake sources, an unresolved issue is whether the S to Lg transfer function is flat at the source. A flat transfer function was empirically proposed by Street *et al.* (1975) and was supported by numerical simulations using flat layered structures by Herrmann & Kijko (1983) and by Campillo *et al.* (1985). Several other authors (eg., Harr *et al.*, 1984, 1986; Mueller & Cranswick, 1985) however, disagree with the idea of a flat transfer function, based on the interpretation of

observational results. They argued that corner frequencies (f_c) estimated using Lg from many intra-plate earthquakes in North America are systematically lower than the f_c values estimated using local S waves from the same earthquakes. Mueller & Cranswick (1985) suggested that the M_0 values estimated using local S waves scaled with f_c^{-3} , whereas M_0 values estimated using the Lg wave scaled with f_c^{-4} . They further suggested that f_c estimated using any regional phases may be biased. Before this study, it seemed unclear to us whether the proposed systematically lower f_c values obtained using Lg were caused by the use of imprecise, path-invariant Q_{Lg} values, or caused by some inherent difference between the excitations of Lg and S waves by sources in realistic 3D media.

Sereno *et al.* (1988), during a novel simultaneous inversion of source spectra and Q using both Lg and Pn waves, assumed path-invariant Q and a scaling of $M_0 \sim f_c^{-3}$ for Pn and Lg excited by both explosions and earthquakes. That scaling disagrees with those proposed by Harr *et al.* (1984, 1986) and Mueller & Cranswick (1985) when earthquakes are concerned. Sereno *et al.* obtained several important and interesting results in their study. One intriguing possibility found in their work is that in terms of seismic moment, the Lg excitation efficiency by explosions appears to be lower by a factor of 0.27 than that by earthquakes. While that result is interesting, we feel that it may have been affected by the various assumptions made, as cautioned by the authors. It would be very interesting to see whether one could obtain the same result without some of the assumptions by Sereno *et al.*

In view of the previous debates/unresolved issues on the Lg excitation, much of this study has been directed to further explore the answers to three fundamental questions in the excitation of the Lg by various seismic sources, summarized in the following:

- (1) Which scaling relation, *i.e.*, $M_0 \sim f_c^{-3}$ or $M_0 \sim f_c^{-4}$, is a better representation of Lg excitation by seismic sources?
- (2) Is the transfer function between local S and Lg flat?
- (3) Is there any fundamental difference between Lg excitation by explosions and by earthquakes?

Scaling between M_0 and f_c values for explosions

The method used in this study (see the beginning of section II) allows us to explore the answers to the questions in the last section without making any *a priori* assumptions of path-invariant Lg Q, or of the scaling between M_0 and f_c . In Figure 8, we plotted the Lg M_0 and f_c for explosion sources, obtained using the explosion source model (equation (1)). The correlation between the logarithm of M_0 and f_c values in Figure 8 is linear and

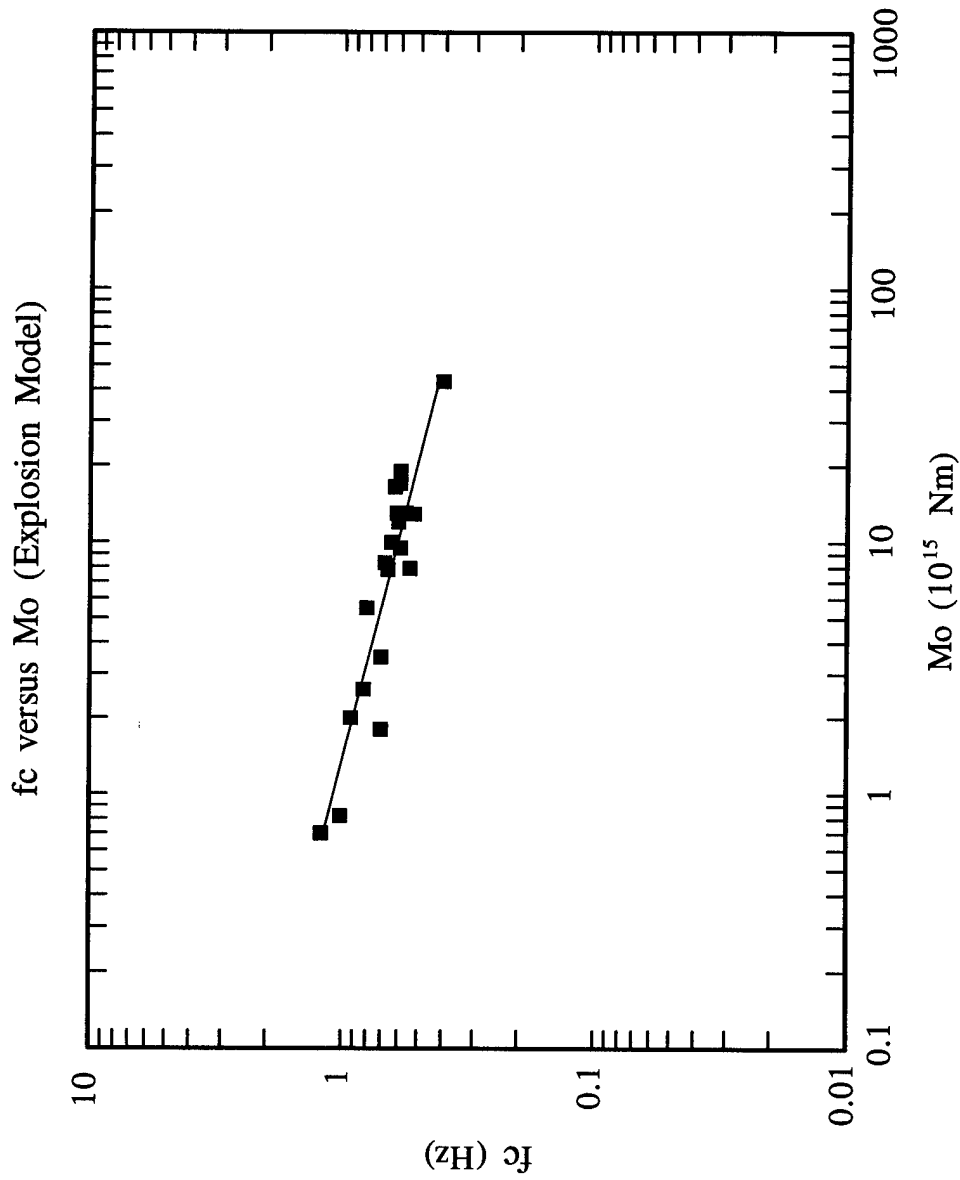


Fig. 8. Logarithm of M_0 (in 10^{15} Nm) versus logarithm of f_c values for the explosions studied, obtained by inverting the Lg spectra using the explosion source model. Straight line represents the linear regression fitting.

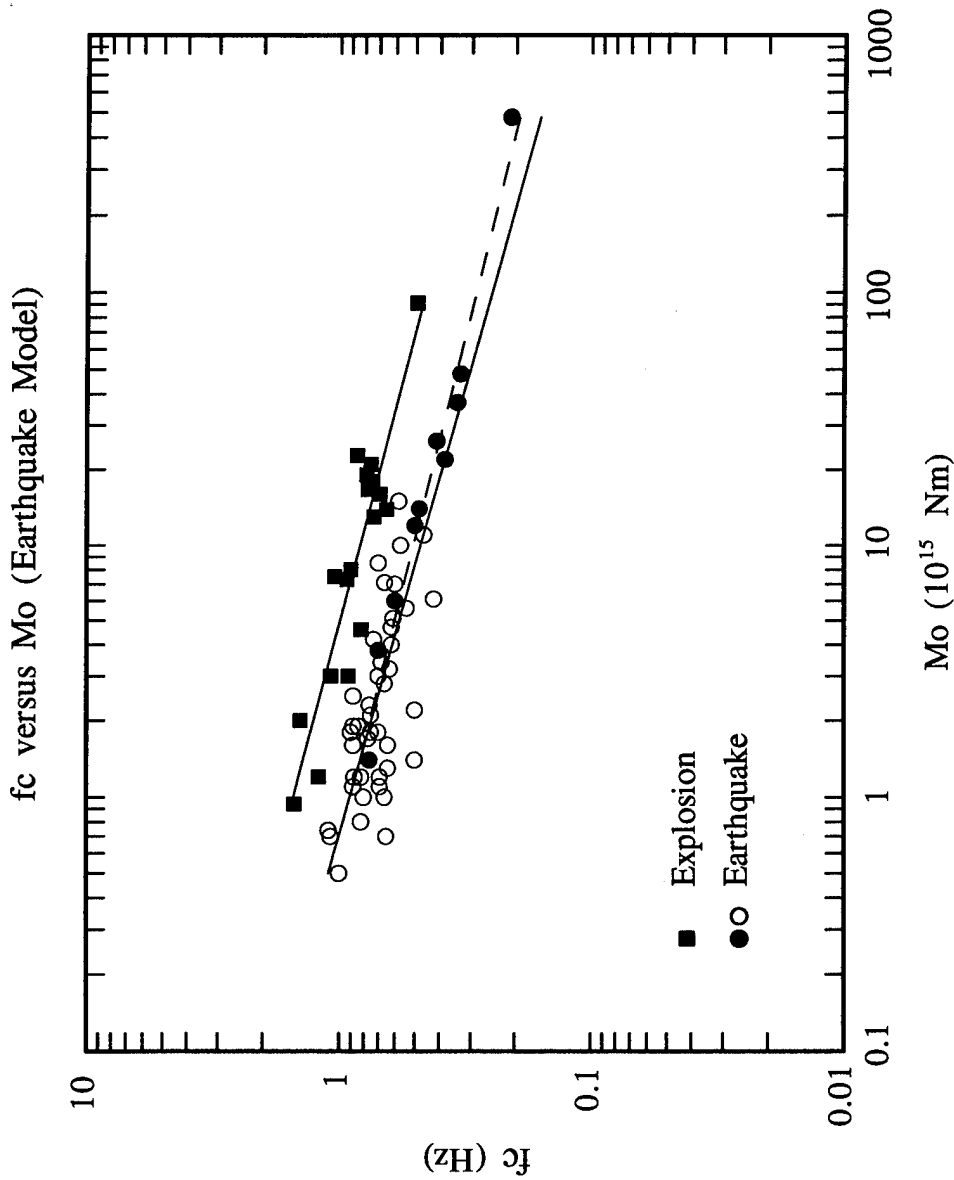


Fig. 9. Logarithm of M_0 versus logarithm of f_c values for the earthquakes and explosions studied, all obtained by using the earthquake source model. Solid circle represent values for earthquakes with Lg recorded at multiple (≥ 3) stations, and open circles represent the values for earthquakes with Lg recorded by fewer stations. Squares represent values for explosions, obtained using the earthquake source model. The two solid lines represent the linear regressions over all earthquakes and all explosions, respectively. The dashed line represent the linear regression over the solid circles only.

a regression analysis yields the following relationship:

$$\log M_0 = 15.12(\pm 0.22) - 3.98(\pm 0.43) \log f_c, \quad (9)$$

which suggests an $M_0 \sim f_c^{-4}$ scaling.

Scaling between M_0 and f_c values for earthquakes

Figure 9 shows the Lg M_0 and f_c values for the earthquake sources, obtained using the earthquake source model (equation (2)). Again, solid circles represent values for earthquakes with Lg recorded by multiple (≥ 3) stations, and open circles represent earthquakes with Lg recorded by only one or two stations. A linear regression over all of the 52 circles yields

$$\log M_0 = 14.85(\pm 0.30) - 3.56(\pm 0.30) \log f_c \quad (10)$$

(solid line in Figure 9). If only solid circles (earthquakes with multiple Lg records) in Figure 9 are used, a linear regression yields

$$\log M_0 = 14.81(\pm 0.16) - 4.04(\pm 0.25) \log f_c \quad (11)$$

(dashed line in Figure 9). The slopes and intercepts in (10) and (11) overlap within the estimated uncertainty, with uncertainties (11) being smaller, probably due to that the corresponding Lg M_0 and f_c values are better constrained when multiple records are available. Both equations suggest that M_0 scale with $f_c^{-\alpha}$, with α being closer to 4 than to 3.

Dependence of M_0 , f_c estimates on theoretical source models

A comparison between the $M_0 - f_c$ scaling derived for the explosions using the explosion source model (Figure 8; equation (9)) with those for the earthquakes using the earthquake model (equations (10) and (11)) will not be useful for the purpose of explosion discrimination since in deriving equation (9), we knew *prior* to the inversions that the events were explosions and accordingly, we have used the explosion source model in the inversions. For the purpose of future discrimination of explosions from earthquakes, it is desirable to simulate a situation when we do not know that the explosions under study are explosions, and treat them as earthquakes in the inversion. Therefore for the 21 explosions, we conducted inversions using the earthquake source model with no overshoot (equation (2) in section II). The resulting M_0 and f_c values are plotted in Figure 9 as squares. A linear regression over these values yields the following relation:

$$\log M_0 = 15.69(\pm 0.23) - 3.83(\pm 0.45) \log f_c. \quad (12)$$

The slope in equation (12) suggests that the scaling of $M_0 \sim f_c^{-4}$ is roughly preserved, even

though an earthquake source model is used in the inversion. On the other hand, the intercept predicted by equation (12) differs from that predicted by (9) by $0.57 (\pm 0.45)$. This means that the M_0 values obtained with the explosion source model are systematically lower by a factor of 0.27 than those obtained with the earthquake source model. This factor coincides with the κ value of 0.27 obtained by Sereno et al. (1988), who studied *different sources (earthquakes and explosion) with the same (explosion) source model*, and proposed that the low κ value indicates a depletion of Lg by explosions. In this study, however, both (12) and (9) are obtained using the same explosion data set, and the factor of 0.27 must have arisen solely because of the model-dependence of the inversion, rather than any fundamental difference between the excitation of Lg by different source types. Physically, this model-dependence is due to the effect of overshoot in the explosion source model (equation (1)), which tends to downsize the seismic moment and/or corner frequency with an observed Lg spectrum. This model dependence makes it more obvious that for the purpose of future explosion discrimination, it is desirable to compare the $M_0 - f_c$ scaling for the explosions with the corresponding scaling for the earthquakes, both derived using the same (preferably earthquake) source model, to look for any systematic differences. In the next section we present such a comparison using results obtained in this study.

Comparison between $M_0 - f_c$ scalings for earthquakes and explosions

As mentioned in the above sections, the circles and squares Figure 9 represent M_0 and f_c values estimated by applying the same (earthquake) source model to earthquakes and explosions, respectively. It is obvious that the distributions of circles and squares in Figure 9 show systematically different trends. At any fixed M_0 level, f_c values for explosions (squares) are systematically higher than those for earthquakes (circles). Consequently, the corresponding linear regression fittings in equations (12) and (10) (or (11)) differ in their intercept by about 0.9, despite that the slopes are all similar (close to -4.0). Physically, this phenomenon means that the Lg excitation by explosions tend to be enriched in high frequency contents, as compared to Lg by earthquakes. We may use this phenomenon to form a basis for discriminanting explosions from earthquakes, but there are two complications associated with this potential discriminant. The first is that some open circles in Figure 9 overlap with the squares, but all of the filled circles are unambiguously separated from the squares. This suggests that a discriminant of explosion using the $M_0 - f_c$ scalings should be used with caution when there are only one or two stations recording the Lg from an seismic event, either due to a sparse-station coverage or a small event size. The second complication is that the phenomenon of an enriched high

frequency Lg excitation by the explosions, observed in this study for the Lop Nor and Balapan Test Sites, appears to be opposite to what was observed for the Nevada Test Site (e.g., Murphy and Bennett, 1982; Bennett and Murphy, 1986; Walter *et al.*, 1995). This suggests that there may be some significant variations in the detailed mechanism of excitation of Lg by explosions with varying tectonic/geological environments, and any discriminant using Lg developed for one test site may not be directly transportable to another test site.

Implications on the difference between excitations of local S and Lg

It has been generally believed that the M_0 and f_c values of earthquakes, estimated using local S waves, tend to scale as $M_0 \sim f_c^{-3}$ (pages 19 and 20). The fact that equations (10) and (11) both suggest slopes that are closer to -4 than -3 suggests that the Lg $M_0 - f_c$ scaling differs from the corresponding scaling derived using local S waves. This probable difference has a profound implication in the mechanism of Lg excitation. The implication, however, is not as direct as desired. A systematic difference between the excitation of local S and Lg waves should be confirmed with more direct comparisons of spectra of S and Lg from the same events.

IV. FUTURE RESEARCH

Based on the results of this study, we recommend that future research be conducted in the following areas:

- (1) Establish more precise, perhaps distance and frequency dependent geometrical spreading for the Lg phase and other regional phases based more substantial observations, made with more data, and synthetics using more realistic velocity structures.
- (2) Explore why the Lg excitation by explosions in different test sites vary in levels of high-frequency contents relative to low-frequency contents.
- (3) Compare frequency content of Lg with that of S waves observed at smaller distances (e.g., the Sg wave), and further explore if and why the local S waves exhibit a different $M_0 - f_c$ scaling than the Lg wave.
- (4) Apply the same methodology in this study to the spectral characteristics of excitation and propagation of other regional phases, particularly the Pg and Pn phases, and systematically evaluate the P/Lg spectral ratio discriminant. Also, test to see if the

P/Lg discriminant is more reliable than the Lg discriminant.

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VII. PROFESSIONAL PERSONNEL

Investigators involved in the two-year research include Dr. Jiakang Xie (the P.I.), Dr. Brian J. Mitchell (the co-P.I.), and Mr. Lianli Cong (a graduate student).

VIII. PAPERS PRESENTED AT MEETINGS

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